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Evolution of Classical 1-D Based Models and Improved Approach for the Characterization of Litz Wire Losses

Asier Arruti, Iosu Aizpuru, Mikel Mazuela, Ziwei Ouyang and Michael A.E. Andersen

Abstract- Prediction of Litz wire losses is a challenging endeavor in the design of high frequency transformers. Although many analytical approaches of varying complexity can be found in literature, classical 1-D based models are the norm in the analysis of wide design spaces. Even then, an important degree of uncertainty exists around the accuracy of the 1-D models. After a critical review of these models and an analysis of their application in Litz wire loss prediction, a new model is proposed which tries to address some of the concerns and limitations of the existing approaches. The models are evaluated against 2-D finite element simulations in many different conditions to ensure the models correctly consider the impact of different parameters. Experimental results are presented to demonstrate the accuracy of the new approach at very high frequencies and high amounts of strands using four prototypes with different Litz wires, designed to correctly match the conditions inherent to any 1-D model.

Index Terms- Litz wire, Copper losses, High-Frequency transformer, Mathematical models.

I. INTRODUCTION

In the design of power electronics converters, an accurate representation of component behavior is critical for the optimal design of the converter. The magnetic devices (transformers and inductors) play a major role in the overall efficiency, power density, weight, and cost of the design [1]-[4], and should be accordingly designed for each specific application.

The inclusion of wide band gap technologies (GaN and SiC transistors) allows to increase the power converter switching frequencies beyond the typical limits of Silicon based IGBT and MOSFET switches [5], [6]. It is well known that the increase in frequency can aid in the reduction of magnetic device size, but this scaling law is limited by the power loss and dissipation capabilities of the devices [1]-[4], [7]. For transformer design, the optimal design frequency of the transformer can range between 10 to 20 kHz for applications in the range of hundreds of kilowatts [8]-[10], to beyond the MHz range for applications up to the kilowatt limit [11]-[14].

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Asier Arruti, Iosu Aizpuru and Mikel Mazuela are with the Faculty of Engineering, Mondragon University, Mondragon, Spain (e-mail: aarruti@mondragon.edu;iaizpuru@mondragon.edu;mmazuela@mondragon.edu).

Ziwei Ouyang and Michael A.E. Andersen are with the Department of Electrical and Photonics Engineering, Technical University of Denmark, Kgs. Lyngby, Denmark (e-mail: ziou@dtu.dk; maea@dtu.dk).

One of the functional constraints in these high frequency transformers is the winding temperature, which becomes complex to evaluate due to the nature of high frequency copper losses. This high frequency losses are primarily due to skin and proximity effect, although in specific cases such us hybrid/integrated magnetics the fringing effect can also be a major contributor [15]-[18]. Litz wire, comprised of transposed insulated strands, can effectively decrease the skin and proximity effects associated with high frequencies, but despite its efficacy in reducing these losses the accurate modelling of Litz wire losses is a challenging endeavor.

A priori, the accurate estimation of Litz wire losses might not appear critical in the examples shown above; lower power applications planar magnetics with PCB windings are used in many cases [11]-[14], but design utilizing Litz wire may still offer significant advantages [6], [19]-[21], or even hybrid PCB-Litz windings might be used [22], [23]. For higher power applications, although the frequencies are much lower than their planar counterparts [8]-[10], due to the requirement of much higher copper cross sections which usually entails higher amounts of strands, the selection of Litz wire has a non-marginal impact in the transformer cost as shown by the cost trends from [24] and [25].

Finite element simulations of the intricate internal geometry are in most cases too computationally heavy for the design process of the transformer, thus analytical approaches are generally preferred. Even though complex analytical models can be found in literature tackling this problem, classical 1-D based approaches, namely Dowell's [26], and Ferreira's [27], [28], methods, are still the norm when it comes to the preliminary evaluation of wide transformer design spaces [4], [8], [10].

Even then, there is still a noticeable uncertainty level around which of these approaches is the best to use, and disagreeing conclusions can be found in literature [27]-[32]. Thus, this paper tries address these concerns while presenting the following main contributions:

– Review of classical 1-D models analysis of the use of these models in literature publications.

– Proposal of 2 alternative models with modifications from Dowell's and Ferreira's models.

– Development of a simulation methodology for fast 2-D FEM simulations for high number of Litz strands, using a more realistic Litz wire internal strand distribution.

– Experimental evaluation of Litz wires losses in very high frequencies.

First, Section II serves as review of these classical 1-D models and their application to the problem of Litz wire losses, while Sections III presents two alternative approaches to model these losses. Comparison of these models against various 2-D finite element simulations of transformer winding configurations are shown in Section IV, where Litz wires of up to 1000 strands are evaluated, and validation against experimental results from transformers adhering to the classical 1-D geometries is included in Section V.

II. Review Of Existing 1-D Models And Their Applications To Litz Wires

For the solution of high frequency losses in copper windings for transformers, two approaches based on the exact mathematical solutions of Maxwell's electromagnetic equations can be found: Dowell's method [26] and Ferreira's method [27], [28]. Note that although both approaches are commonly named after the respective paper's authors, the mathematical solutions of these problems can be dated back to previous works: solutions for rectangular windings resembling those presented by Dowell can be found in 1940 [33], and solutions for cylindrical windings close to Ferreira's formulations are found in 1922 [34]. Still, for the shake of simplicity and familiarity, the approaches will be classified into models based on Dowell's approach (solution for rectangular wires) and models based on Ferreira's approach (solution for circular wires).

Both approaches use an equivalent high frequency resistance to define the power losses, P_{loss}, as a function of effective current,

$$P_{\text{loss}} = I_{\text{rms}}^2 \cdot (R_{\text{DC}} \cdot F_{\text{R}}) \tag{1}$$

where the equivalent high frequency resistance is the DC resistance, R_{DC}, multiplied by the resistance ratio F_{R}.

Dowell's approach:

The Dowell's solution for F_{R} is obtained when Maxwell's electromagnetic equations are solved for an infinitely tall rectangular conductor under the assumption of unidimensional field (in the y direction), as shown in Fig. 1a. This situation is equivalent to finite rectangular conductors placed between two high permeability mediums, which is a common case in foil transformer windings. In these circumstances, according to Dowell's work [26], the average F_{R} value for the windings takes the form

$$F_{\text{R}} = \Delta\left(\frac{\sinh 2\Delta+\sin 2\Delta}{\cosh 2\Delta-\cos 2\Delta} + \frac{2(m^2-1)}{3}\frac{\sinh \Delta-\sin \Delta}{\cosh \Delta+\cos \Delta}\right) \tag{2}$$

although other equivalent forms can be found in [33] and [35]. Here, m is the number of layers (number of conductors in the x direction) and Δ is the penetration depth defined as a function of the skin depth δ

$$\Delta = \frac{w_{\text{c}}}{\delta} \tag{3}$$

$$\delta = \sqrt{\pi f \sigma \mu}^{-1} \tag{4}$$

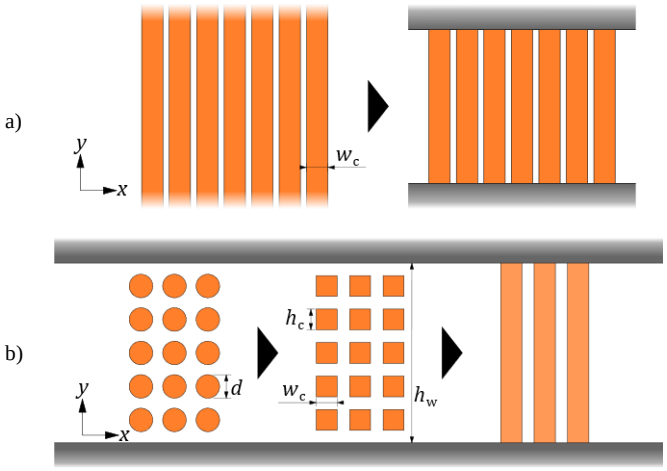

Fig. 1: a) Winding structure for the Dowell approach and b) application to cylindrical conductors.

for a given frequency f, conductivity σ and permeability μ. In the case of transformers, the winding area permeability (air) is used, so that $\mu = 4\pi \cdot 10^{-7}$.

The case presented in [26] assumes the conductors occupy the full window area height, but this is rarely the case in real transformers, since they are usually made of various conductors with gaps in the y direction. The 1-D magnetic field in the borders of a layer of N vertically stacked conductors are equal to a solid layer carrying N times the current, thus Dowell considers the F_{R} of a layer of N vertically stacked rectangular conductors is the same of a solid layer of equivalent conductivity, Fig. 1b. Then, the skin depth is modified so that

$$\delta' = \sqrt{\eta}^{-1}\delta \tag{5}$$

$$\eta = N\frac{h_{\text{c}}}{h_{\text{w}}} \tag{6}$$

where δ' is the 'effective skin depth', and η is the ratio between the height occupied by the conductors $N \cdot h_{\text{c}}$, to the height of the winding area h_{w}, referred to as 'conductor spacing factor', 'layer porosity', 'layer copper factor', or most commonly 'porosity factor'. In the case of cylindrical wires of diameter d, Dowell first transforms these into rectangular conductors of equal cross section, so that the height and width of the conductors are

$$h_{\text{c}} = w_{\text{c}} = \sqrt{\frac{\pi}{4}}d \tag{7}$$

and then solves the problem using the proper porosity factor.

Note that the transformation from circular to square conductors using the same cross-section is not supported by any physical meaning [27], [36], and the relation of $w_{\text{c}} = (3\pi/16)^{1/4}d$ has been found to give better results [37]. Also, several works have pointed out that the porosity factor entails an error in Dowell's calculations, since skin depth as a physical parameter cannot be influenced by the porosity factor [27], and Ampere's law is mismatch for $\eta \neq 1$ [36], thus Dowell's approach is considered only accurate for high porosity factors.

Ferreira's approach:

To evaluate high frequency copper losses with circular conductors circumventing the intrinsic problems from the Dowell's approach, Ferreira presented another method using the mathematical solution in cylindrical coordinates based on Bessel functions [27], [28]. To do so, Ferreira's work uses the orthogonality principle between skin and proximity losses, postulated in [38], meaning that both loss phenomena can be calculated independently and added together. This can be demonstrated from Dowell's approach, where (2) can be rewritten as

$$F_{\mathrm{R},m} = \frac{\Delta}{2}\left(\frac{\sinh\Delta+\sin\Delta}{\cosh\Delta-\cos\Delta} + (2m-1)^2\frac{\sinh\Delta-\sin\Delta}{\cosh\Delta+\cos\Delta}\right) \tag{8}$$

where the $F_{\mathrm{R},m}$ is defined per layer instead of the average F_{R} of the winding section. Here, the first term is equal to the solution of skin effect for square conductors, thus according to the orthogonality the second term defines the proximity losses in the m^{th} layer. As demonstrated by [35], the term $(2m-1)$ relates to the magnetic field strength applied to each layer, so it can be concluded that the losses grow as a function of the square of this applied magnetic field.

In this case, the losses for circular conductors can be obtained by first solving the skin and proximity losses independently and adding them together. The resulting expression for circular wires is

$$F_{\mathrm{R},m} = \frac{\gamma}{2}\left(\frac{\mathrm{ber}\,\gamma\,\mathrm{bei}'\,\gamma-\mathrm{bei}\,\gamma\,\mathrm{ber}'\gamma}{\mathrm{ber}'^2\,\gamma+\mathrm{bei}'^2\,\gamma} - 2\pi(2m-1)^2\frac{\mathrm{ber}_2\,\gamma\,\mathrm{ber}'\,\gamma+\mathrm{bei}_2\,\gamma\,\mathrm{bei}'\gamma}{\mathrm{ber}^2\,\gamma+\mathrm{bei}^2\,\gamma}\right), \tag{9}$$

where ber and bei are Kelvin functions, representing the real and imaginary parts of the Bessel function respectively, and in this case

$$\gamma = \frac{d}{\delta\sqrt{2}} \tag{10}$$

is the adapted definition for the penetration depth in round conductors.

The expression achieved in [27] is based on that of [26] with modified functions, since Ferreira simply replaces the skin effect and proximity effect expressions with the ones for circular wires. This is purposely done to compare the discrepancies in the F_{R} values with Dowell's results. This does not mean that (9) can be used as is, since Ferreira clearly expresses the limitations of the porosity factor proposed by Dowell. Based on Ferreira's demonstration, the more adequate definition for high frequency square conductor losses taking the orthogonality principle is

$$F_{\mathrm{R},m} = S_{\mathrm{R}} + G_{\mathrm{R}}H_{\mathrm{e}}^2 \tag{11}$$

$$S_{\mathrm{R}} = \frac{\Delta}{2}\frac{\sinh\Delta+\sin\Delta}{\cosh\Delta-\cos\Delta} \tag{12}$$

$$G_{\mathrm{R}} = \frac{\Delta}{2}\frac{\sinh\Delta-\sin\Delta}{\cosh\Delta+\cos\Delta} \tag{13}$$

where S_{R} and G_{R} are the skin and proximity factors respectively, and H_{e} is the external magnetic field applied to the conductor. The analogue factors for round conductors are then

$$S_{\mathrm{R}} = \frac{\gamma}{2}\frac{\mathrm{ber}\,\gamma\,\mathrm{bei}'\,\gamma-\mathrm{bei}\,\gamma\,\mathrm{ber}'\gamma}{\mathrm{ber}'^2\,\gamma+\mathrm{bei}'^2\,\gamma} \tag{14}$$

$$G_{\mathrm{R}} = -\gamma\pi\frac{\mathrm{ber}_2\,\gamma\,\mathrm{ber}'\,\gamma+\mathrm{bei}_2\,\gamma\,\mathrm{bei}'\gamma}{\mathrm{ber}^2\,\gamma+\mathrm{bei}^2\,\gamma} \tag{15}$$

The porosity factor also relates the field generated by a solid conductor occupying the full winding area height and the magnetic field generated by conductors not occupying the full height, as illustrated in Fig. 1b. Thus, taking this reduction of the magnetic field

$$F_{\mathrm{R},m} = \frac{\Delta}{2}\left(\frac{\sinh\Delta+\sin\Delta}{\cosh\Delta-\cos\Delta} + \eta^2(2m-1)^2\frac{\sinh\Delta-\sin\Delta}{\cosh\Delta+\cos\Delta}\right) \tag{16}$$

should result in a more accurate description of high frequency losses for square conductors than (8), where Δ is calculated with the δ instead of δ'. It follows that the same logic must be applied to (9) for cases where $\eta \neq 1$. Ferreira does not present such an equation, and instead proposes to evaluate the magnetic field generated by the rest of the wires and use this in the calculations avoiding the necessity of η, resulting in a more generalized solution. An expression based in (9) with the inclusion of the porosity factor can be found in [39], [40], taking the form

$$F_{\mathrm{R},m} = \frac{\gamma}{2}\left(\frac{\mathrm{ber}\,\gamma\,\mathrm{bei}'\,\gamma-\mathrm{bei}\,\gamma\,\mathrm{ber}'\gamma}{\mathrm{ber}'^2\,\gamma+\mathrm{bei}'^2\,\gamma} - 2\pi\eta^2(2m-1)^2\frac{\mathrm{ber}_2\,\gamma\,\mathrm{ber}'\,\gamma+\mathrm{bei}_2\,\gamma\,\mathrm{bei}'\gamma}{\mathrm{ber}^2\,\gamma+\mathrm{bei}^2\,\gamma}\right), \tag{17}$$

but is limited to transformers with winding layers composed of the same number of round conductors.

Remark that in the Dowell's based approximation with conductors filling the complete vertical height, the $(2m-1)$ term from (8) related to H_{e} is present even in single layer windings. Using the 1-D field generated by all the wires in the evaluation of H_{e} instead of the rest of the wires results in a better mathematical implementation of orthogonality principle for typical transformer configurations.

Although Ferreira's approach is based on the Bessel function solutions for circular wires, many authors have reported better accuracy using Dowell's approach combined with the equivalent cross-section transformation (7), [30], [32], [36]. Note that the comparison from [32] might not be the most reliable due to some flaws detected in the analysis:

– The *Dowell Method (Dowell 1)* and *Modified Dowell Method (Dowell 2)* models, equivalent to (2), are presented as different models when the first is a specific case from the second.

– The *Ferreira Method (Ferreira 1)* model, analogue to (14), not only applies the porosity factor to reduce the external field, but also uses the δ', thus overcorrecting the losses due to proximity effect.

– The *Original Ferreira Method (Ferreira 2)* and *Modified Ferreira Method (Reatti)* models, equivalent to (9) and (15) are presented as different models when the first is a specific case from the second as demonstrated above.

– The ratio between the *Original Ferreira Method (Ferreira 2)* and *Modified Ferreira Method (Reatti)* model power losses, (9) and (15), cannot differ by a factor lower than η^2 (assuming 100 % proximity losses), while the

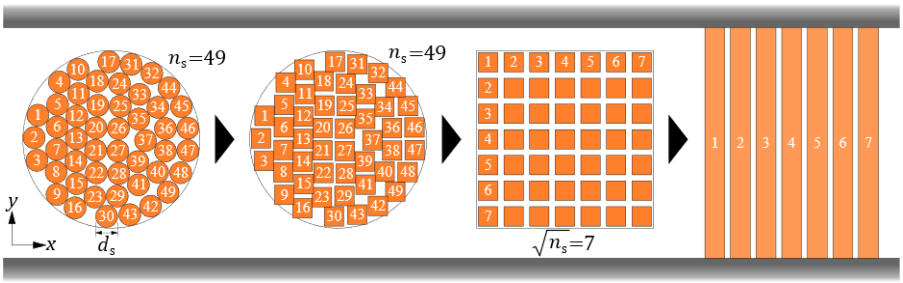

Fig. 2: Transformation of Litz wire of n_s circular strands into rectangular strands and rearranged into an $\sqrt{n_s}$ by $\sqrt{n_s}$ matrix to apply Dowell's approach.

reported results appear to go beyond this limit. Note that the format of the presented results (3-D contour maps) complicates the extraction of precise data to prove this point.

– The results displayed differ noticeably from the reported in [30], where typical errors of ~5 % and ~1 % are presented for Dowell's and Ferreira's models respectively in the low frequency regions.

Litz loss modelling with 1-D approaches:

The application of the 1-D approaches to Litz wire conductors encompasses new challenges for the geometrical definition of the problems. To do so, two approaches can be identified in the literature: transformation of the Litz wire into equivalent winding layers, and the addition of the internal proximity factor to the losses.

Dowell's equation on Litz wires (Wojda's approach):

The concept behind the first approach is illustrated in Fig. 2, where each Litz wire bundle composed of n_s number of strands is transformed into a square wire matrix of $\sqrt{n_s}$ by $\sqrt{n_s}$ strands [41]. Then it follows that the porosity factor (6) is modified as

$$\eta = N\sqrt{n_s}\sqrt{\frac{\pi}{4}}d_s/h_w \tag{18}$$

where $\sqrt{\pi/4}\,d_s$ is the equivalent square strand width of each circular strand using the cross-sectional transformation (7).

Then the problem can be solved using (8) using where

$$\Delta = \sqrt{\frac{\pi}{4}}d_s/\delta\sqrt{\eta} \tag{19}$$

is the penetration ratio according to [26] or using (15) where

$$\Delta = \sqrt{\frac{\pi}{4}}d_s/\delta \tag{20}$$

is the penetration ratio for the external magnetic field correction method proposed in [27].

Geometrical transformation also means that the number of winding layers has increased by a factor $\sqrt{n_s}$. Since the definition of layer losses physical meaning in this case, instead of obtaining the equivalent $F_{R,m}$ for a specific layer the average F_R of all winding layers must be used to solve the problem. Transforming (8) and (15) back into an analogue form of (2) results in

$$F_R = \frac{\Delta}{2}\left(\frac{\sinh\Delta+\sin\Delta}{\cosh\Delta-\cos\Delta}+\left(\frac{4(n_s\cdot m^2-1)}{3}+1\right)\frac{\sinh\Delta-\sin\Delta}{\cosh\Delta+\cos\Delta}\right) \tag{21}$$

and using Ferreira's modified Dowell's equation with unmodified skin effect

$$F_R = \frac{\Delta}{2}\left(\frac{\sinh\Delta+\sin\Delta}{\cosh\Delta-\cos\Delta}+\eta^2\frac{2(n_s\cdot m^2-1)}{3}\frac{\sinh\Delta-\sin\Delta}{\cosh\Delta+\cos\Delta}\right). \tag{22}$$

This approach was popularized by [31], thus is commonly referred to as Wojda's approach, although as explained before the concept of transformation of Litz strands into an square grid was presented previously by Alex Van Den Bossche and Vencislav Cekov Calchev in 2005 (equation 2.23 in [41]).

$$F_R = \Delta\cdot\left(\frac{\sinh 2\Delta+\sin 2\Delta}{\cosh 2\Delta-\cos 2\Delta}+0.95\frac{2(n_s\cdot m^2-1)}{3}\frac{\sinh\Delta-\sin\Delta}{\cosh\Delta+\cos\Delta}\right) \tag{23}$$

Here some minor modifications and correction factors are added to improve the accuracy of the results based:

– The formulations for the penetration depth and porosity factor are defined as

$$\Delta = \left(\frac{\pi}{4}\right)^{3/4}\frac{d_s}{\delta}\sqrt{\eta} \tag{24}$$

$$\eta = \sqrt{n_s}N\frac{d_s}{h_w} \tag{25}$$

where instead of using the strategy originally described by Dowell [26], the porosity factor is defined using the diameter of the circular conductor. Mathematically both approaches are equivalent, but the definition (6) losses physical sense.

– The proximity effect losses are decreased by a factor of 0.95 (π/3) according to the findings reported by [41]. Unlike in expression (8), the proximity losses in expression (2) are not solely represented by the second term (first term combines eddy losses and part of the proximity losses as demonstrated by Ferreira [27]), thus a small error is introduced in the implementation of this correction factor. Still, the error introduced is negligible in most practical cases with Litz wires.

Ferreira's approach for Litz wires:

In [42] Ferreira proposes another approach to model Litz wire high frequency losses. Instead of reordering the Litz wire strands into a matrix, the external magnetic field applied to the Litz wire (average H_e value, or H_e value in the center of the Litz wire bundle) and the internal collective field of the stranded conductors (H_i) are used, as illustrated on Fig. 3.

The Litz strand power losses due to the internal magnetic field are

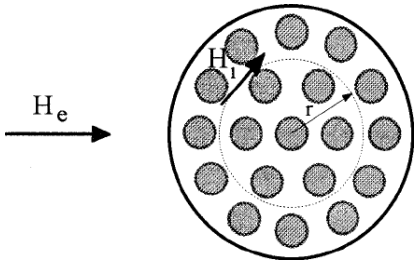

Fig. 3: External magnetic field H_e generated by the rest of the wires and the internal magnetic field H_i generated by the strands. [39]

$$P_i = \frac{G_R I_b^2}{8\pi^2 r_b^2 n_s} \tag{26}$$

where I_b is total current in the Litz bundle and r_b is the radius of the Litz bundle [42]. Then, from the demonstration of orthogonality between H_e and H_i, it follows that the high frequency losses are,

$$F_{R,m} = S_R I^2 + G_R(H_e^2 + H_i^2) \tag{27}$$

and H_i is derived from (22)

$$H_i = \frac{I_b}{\sqrt{2}\cdot 2\pi\cdot r_b\sqrt{n_s}}. \tag{28}$$

It follows that the γ parameter is defined using the Litz wire strand diameter, not the Litz wire bundle diameter.

A similar approach is presented in [39], where the losses due to the external magnetic field H_e and internal magnetic field H_i are taken separately, but the power loss definitions are normalized as functions of layer m instead of magnetic fields (porosity factors mathematically defined). Unfortunately, the approach from [39] has some caveats compared to [42]:

– The skin effect factor S_R is divided by the number of strands, thus at very low frequencies, where $\gamma \ll d_s$ and proximity losses are negligible ($G_R \approx 0$), F_R becomes $1/n_s$.

– The internal Litz field losses are multiplied by the layer term, meaning that these increase as a function of the number of layers. This goes against the orthogonality principle between external and internal Litz wire losses described in [42].

– The definition of an external porosity factor η_1 for the external field and an internal porosity factor η_2 for the internal field are mentioned, which make physical sense, but only η_1 can be described, as the Litz wire strands are distributed in 2-D and η describes a 1-D distribution. According to [39], turn-to-turn, layer-to-layer, and strand-to-strand distances are to consider when defining η_2, as well as the non-uniformity of the magnetic field, but not further explanations of how to do so are given.

Despite the different attempts to implement these equations, the approach from [39] did not yield satisfactory results, with error orders of magnitudes higher than the rest of the models tested in Section III.

III. Proposed Models

In Ferreira's approach for Litz wire models, two sources of proximity losses are assumed, proximity losses generated by the external magnetic field H_e and proximity losses from the internal collective magnetic field H_i. The problem with this assumption is that it does not accurately depict the classical transformer geometry where the 1-D field analysis is valid. When the Litz wire is structured in layers inside the core window, the 1-D magnetic field across the x dimension is not constant, taking the form of sigmoid due to the circular geometry of the Litz wire. Since the proximity losses grow quadratically with H_e, this average field results in an underestimation of the losses.

For the transformer winding structure in interest, alternatively the classical 1-D field assumption can be used. Note that this field distribution accounts for the magnetic field generated by the strands as part of the problem, thus the inclusion of the H_i losses is not required. This logic follows the remark made above, where the H_e parameter should not only define the field generated by the rest of the conductors, but also the conductor under analysis itself.

Then, it follows that (15) can be used to evaluate the losses in the strands, but the complex definition for $\eta_{(x)}$ and $m_{(x)}$ along the variable geometry in the x direction makes this complicated. To circumvent this problem, the magnetic field related physical definition of layer presented in [35] can be used, where a layer is defined as

$$m = \frac{H_{(x=w_c)} - H_{(x=0)}}{H_{(x=w_c)}}. \tag{29}$$

where $H_{(x=0)}$ and $H_{(x=w_c)}$ are the magnetic fields to the left and right of the conductor. In the case of circular conductors w_c is replaced by the strand diameter d_s.

Assuming that the total current in the Litz wire bundle is n_s amperes, so that the current in each strand is normalized to 1 A, the effective increase in magnetic field per strand is

$$H_{(x=d_s)} - H_{(x=0)} = \frac{N_b}{h_w}, \tag{30}$$

where N_b is the amount of Litz wire bundles stacked vertically per layer, and h_w is the height of the core window. Then, the Litz wire geometry can be simplified, so that the n^{th} strand in the x direction has a 1-D average equivalent magnetic field H_e of

$$H_{e(n)} = \frac{N_b}{h_w}\frac{(2n-1)}{2}. \tag{31}$$

for m Litz wire layers, then $n = [1..m\cdot n_s]$. Lastly, since the losses of each individual strand using Ferreira's equation for round conductors are

$$F_{R,n} = \frac{\gamma}{2}\left(\frac{\mathrm{ber}\,\gamma\,\mathrm{bei}'\,\gamma - \mathrm{bei}\,\gamma\,\mathrm{ber}'\gamma}{\mathrm{ber}'^2\,\gamma + \mathrm{bei}'^2\,\gamma} - \frac{N_b}{h_w}\frac{(2n-1)^2}{2}\frac{\mathrm{ber}_2\,\gamma\,\mathrm{ber}'\,\gamma + \mathrm{bei}_2\,\gamma\,\mathrm{bei}'\gamma}{\mathrm{ber}^2\,\gamma + \mathrm{bei}^2\,\gamma}\right), \tag{32}$$

the definition for the average high frequency loss increment in all strands becomes

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$$F_{\mathrm{R}} = \frac{\gamma}{2}\left(\frac{\mathrm{ber}\,\gamma\,\mathrm{bei}'\,\gamma - \mathrm{bei}\,\gamma\,\mathrm{ber}'\gamma}{\mathrm{ber}'^2\,\gamma + \mathrm{bei}'^2\,\gamma} - \left(\frac{N_{\mathrm{b}}}{h_{\mathrm{w}}}\right)^2\left[\frac{1}{3}(m^2 n_{\mathrm{s}}^2 - 1) + \frac{1}{4}\right]\frac{\mathrm{ber}_2\,\gamma\,\mathrm{ber}'\,\gamma + \mathrm{bei}_2\,\gamma\,\mathrm{bei}'\gamma}{\mathrm{ber}^2\,\gamma + \mathrm{bei}^2\,\gamma}\right). \quad (33)$$

Expression (32) also allows to approximate the losses per Litz strand ordered in ascending x direction, which allows for a more in-depth comparison with FEM simulations than (33). Unlike the classical Ferreira's approach (27), the analysis of the field per strand removes the intrinsic field averaging errors, which are more noticeable at low number of layers.

Alternatively, based on FEM simulation results, a modified version of Wojda's approach is also proposed. The methodology is the same, but the key difference is that instead of representing the Litz wire strands in an square grid of $n_{\mathrm{s}}^{0.5}$ by $n_{\mathrm{s}}^{0.5}$ dimensions, the strands are distributed into a rectangular grid of $n_{\mathrm{s}}^{0.5-\epsilon}$ stacks and $n_{\mathrm{s}}^{0.5+\epsilon}$ layers, maintaining the total amount of strands equal ($n_{\mathrm{s}}^{0.5-\epsilon} \cdot n_{\mathrm{s}}^{0.5+\epsilon} = n_{\mathrm{s}}$). The value of ϵ is adjusted from the simulation results, where $\epsilon = 0.05$ has been found to give the best overall results. The proximity loss correction factor has also been modified, instead of 0.95 the value $3/\pi$ is used according to [41]. These modifications result in the expression

$$F_{\mathrm{R}} = \Delta \cdot \left(\frac{\sinh 2\Delta + \sin 2\Delta}{\cosh 2\Delta - \cos 2\Delta} + \frac{\pi}{4}\frac{2(n_{\mathrm{s}}^{1+2\epsilon} \cdot m^2 - 1)}{3}\frac{\sinh \Delta - \sin \Delta}{\cosh \Delta + \cos \Delta}\right). \quad (34)$$

$$\Delta = \sqrt{\frac{\pi}{4}\frac{d_{\mathrm{s}}}{\delta}}\sqrt{\eta} \quad (35)$$

$$\eta = n_{\mathrm{s}}^{0.5-\epsilon} N \sqrt{\frac{\pi}{4}\frac{d_{\mathrm{s}}}{h_{\mathrm{w}}}} \quad (36)$$

IV. Comparison With Finite Element Simulations

To evaluate the accuracy of the different models, 2-D simulations of Litz wires forming transformer winding structures are used. The simulations are carried using the Finite Element Method Magnetics (FEMM) software [43].

To model the current distribution between Litz wire strands, an equal current is defined in all strands, simulating an ideal strand transposition [44]. According to [37], [45], [46], in wires with hundreds of strands, the deviation from homogeneity on strand level has only a very minor influence of bundle level effects.

The assumption of equally distributed current in the strands also aids in the simulation of Litz wires with high numbers of strands, since at macroscopic level the Litz wire can be replaced as a solid wire carrying a homogeneous current density. Assuming a normalized current of 1 A per strands, then

$$J_{\mathrm{b}} = \frac{4n_{\mathrm{s}}}{\pi d_{\mathrm{b}}^2} \quad (37)$$

where J_{b} is the current density in an equivalent solid wire of equal diameter to the Litz wire bundle d_{b}. This simplification effectively reduces the computational cost of the FEM simulation, allowing to model Litz wires of high amounts of strands more efficiently.

To model the geometrical strand distribution inside the Litz wire, different assumptions have been used in literature; radially distributed layers [39], square packing [47] and triangular packing [44]. In reality, the transposition of the strands results in a more disordered arrangement, thus, an alternative distribution is proposed in this paper based on the best-known circles in a circle configuration [48]. The angle of each Litz wire is rotated randomly to better depict the unorganized distribution of strands. The differences between internal strand distributions are illustrated in Fig. 4.

The FEM simulations evaluate the impact of four key geometrical parameters in the accuracy of the different 1-D models; number of layers, number of strands, core window fill factor, and Litz wire packing factor. For the base case, a configuration of 3 layers, 200 strands, layer fill factor of $\pi/4$ (closely stacked Litz bundles) and packing factor of 0.65 is considered. The impact of the different parameters is evaluated by deviating from the base case.

The results for the baseline case are shown in Fig. 5. Here the Dowell on Litz (21), Dowell on Litz with Ferreira's modification (22) and Wojda's method (23) are shown for the Dowell based approaches, as well as Ferreira's Litz wire model (27). The proposed approaches based on Wojda's method (34) and Ferreira's method (33) are also displayed. As can be seen, both (21) and (22) overestimate the losses by a factor of $\pi/3$ as reported by [41], although the modification introduced by Ferreira is noticeable worse at predicting very high frequency losses. Wojda's approach takes this overestimation into account, achieving very accurate predictions, but starts underestimating the losses past $d_{\mathrm{s}}/\delta = 1$. On the contrary, Ferreira's approach for Litz wires is slightly less accurate at low frequencies but retains a higher accuracy at very high frequencies ($d_{\mathrm{s}}/\delta > 1$). Lastly, the proposed method based on Wojda's approach can accurately estimate the losses up to higher frequencies than the original Wojda's model, while the

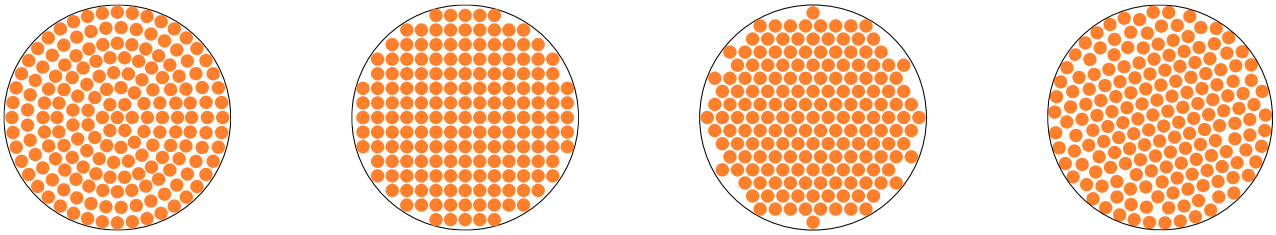

Fig. 4: Different Litz wire strand distribution modelling approaches: a) radial layers used in [39], b) square packing used in [47], c) triangular packing used in [44], and d) best-known circles in circle with random rotation (this work).

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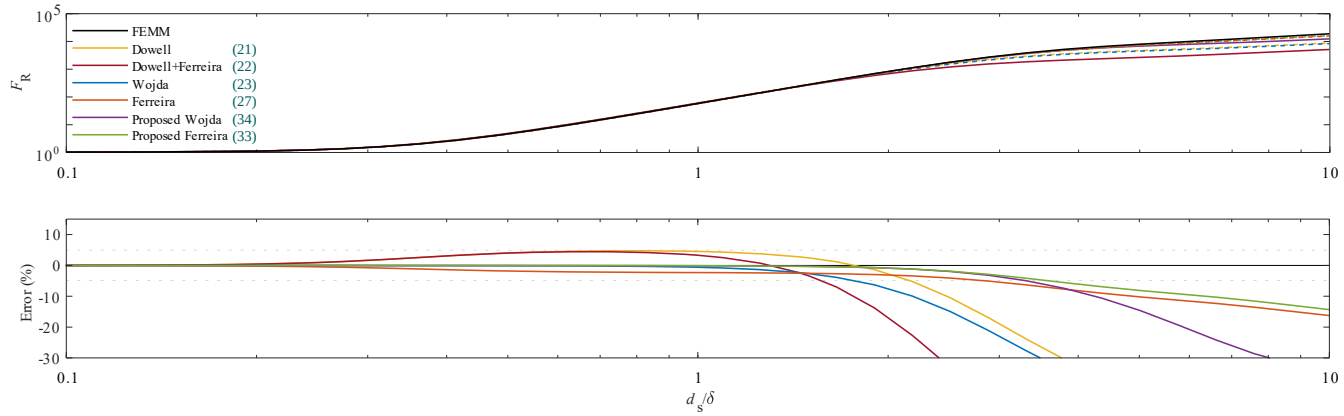

Fig. 5: High frequency winding loss predictions of the different models in the baseline case ($m = 3$, $n_s = 200$, PF $= n_s d_s^2/d_b^2 = 0.65$, $\eta = \pi/4$).

proposed approach based on Ferreira's equations achieve similar results to Ferreira's method but with higher accuracies. Based on these results, going forward only Wojda's model, Ferreira's model for Litz wires, the proposed modified Wojda's approach and the proposed model based on Ferreira's approach will be analyzed.

When the number of layers is modified, the results change noticeably as can be seen in Fig. 6. Wojda's model appears to be quite insensitive to the change of layers, while Ferreira's model becomes more accurate as the number of layer's is increased since the impact of the error introduced by the averaged H_e diminishes. The opposite is also true for low amounts of layers, where Ferreira's model heavily underestimates the proximity losses for the results with 1 layer. Both the proposed models show similar evolutions with the change in number of layers, they retain overall higher accuracies than the preceding models, but overestimate the losses by up to 10-15 % at extremely high frequencies, although these are still better predictions than their counterparts at extreme frequencies.

On the opposite, as shown by Fig. 7, all methods appear to be quite immune to the change in the number of strands, with very similar results to the baseline case. Wojda's method appears to barely change with the number of strands, showing discrepancies of less than 0.005 % between the errors for 12 and 1000 strand simulations, while both Ferreira's model and the proposed Ferreira's model only start changing at very high frequencies ($d_s/\delta > 5$). The only model noticeably affected is the proposed modified Wojda's approach, showing improved accuracies as the number of strands increases. It should be noted that it starts showing a slight overestimation of losses in the transition from high frequency to very high frequency losses ($d_s/\delta \sim 3$), like the results with low amounts of layers. FEM simulations with more than 1000 strands are deemed too computationally expensive to evaluate, thus, although the estimations are better than Wojda's approach, the accuracy for very high amounts of strands is uncertain.

The packing factor has a similar impact in all the models, as demonstrated in Fig. 8. The accuracy of Wojda's models is slightly increased for low packing factors, although the model is still unable to give accurate predictions at very high frequencies. Both Ferreira's and the proposed Ferreira's models do not change at all until very high frequencies. The proposed modified Wojda's method shows a similar evolution than the

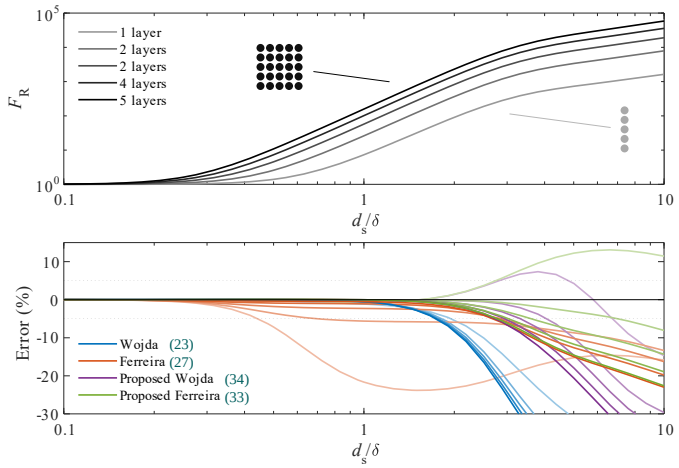

Fig. 6: High frequency winding loss predictions of the different models for different number of layers (m).

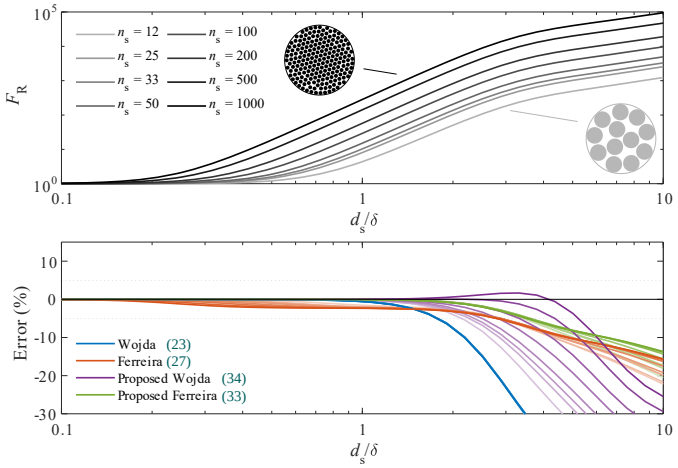

Fig. 7: High frequency winding loss predictions of the different models for different number of Litz strands (n_s).

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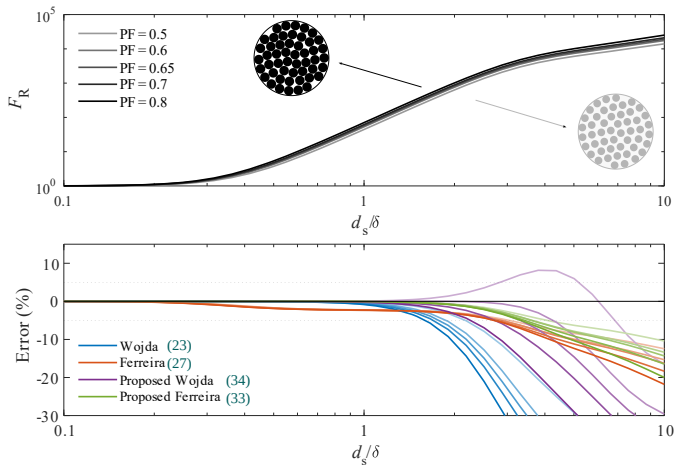

Fig. 8: High frequency winding loss predictions of the different models for different Litz bundle packing factors (PF = $n_s d_s^2/d_b^2$).

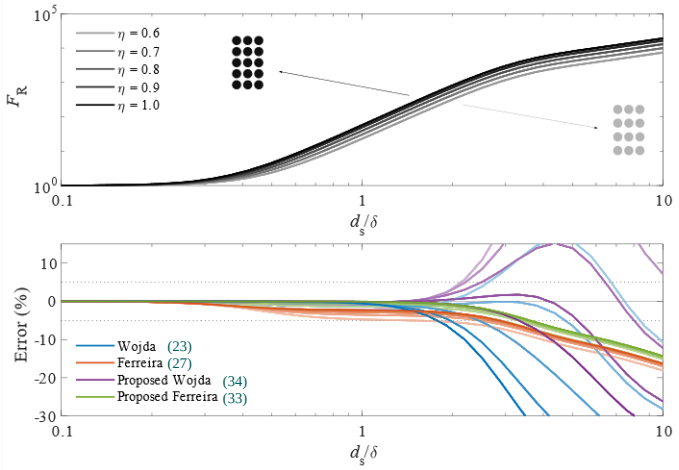

Fig. 9: High frequency winding loss predictions of the different models for different porosity factors (η).

original Wojda's model, but with a much higher sensibility, overestimating the losses in the high frequency to very high frequency transition at low packing factors.

Lastly, the impact of the vertical spacing between wires (porosity factor) is shown in Fig. 9. Both Dowell based approaches change noticeably with the porosity factor, the proposed modified Wojda's approach being better at high porosity factors but showing a higher loss overestimation at very high frequencies. On the contrary, both Ferreira based approaches appear to be immune to these changes. Although hard to see in Fig. 8, the accuracy of all the models changes similarly up to $d_s/\delta = 1$, all showing around –2.5 % error at low porosity factors compared to the baseline case.

To quantify the results, TABLE I resumes the obtained maximum and minimum errors of the different methods in the different simulations at d_s/δ = 1. Also, the d_s/δ values where the error surpasses 5 % are summarized in TABLE II. The colors are used to rank the performances of the models.

As mentioned before, the proposed Ferreira's approach-based method has a unique advantage, the capability to estimate the losses per strand individually using the alternative expression (32). To show the accuracy of this estimation of losses per strand, the losses per strand of the baseline case at $d_s/\delta = 1$ and the estimated losses per strand from (32) are shown in Fig. 10. Here, the x axis represents the horizontal position of the individual strands. Although the strands in the border of the Litz wire show a higher error than the rest (rightmost strand has a prediction error of –1.275 %), due to the circular geometry of the Litz wires, more strands are found in the center than in the borders, thus the impact of this error is reduced.

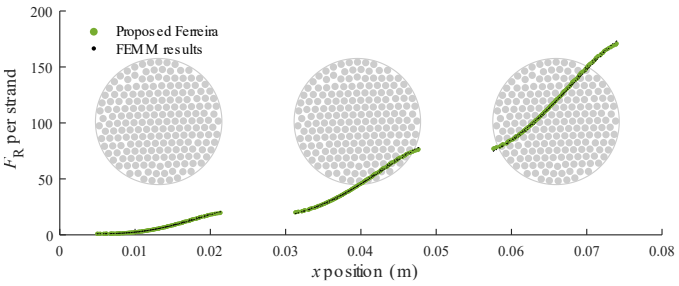

Fig. 10: High frequency winding loss per strand, simulation results and prediction using the proposed Ferreira's approach per strand (32).

TABLE I: ERRORS (%) AT d_s/δ = 1 (LOWER |E| = BETTER)

Simulation	Wojda	Ferreira	Modified Wojda	Proposed Ferreira [REF]
Layers	–0.5063 > E > **–1.1354**	–0.4958 > E > **–21.652**	–0.0159 > E > **–0.6862**	–0.0159 > E > **–0.6315**
Strands	–0.4202 > E > **–0.6147**	–1.5418 > E > –2.3488	+0.0204 > E > **–0.3104**	+0.0848 > E > **–0.1133**
Packing	–0.3090 > E > **–0.8344**	–2.2473 > E > **–2.3545**	+0.0659 > E > **–0.2368**	–0.0262 > E > **–0.1297**
Porosity	–0.5677 > E > **–2.4147**	–2.2994 > E > **–4.7313**	–0.0809 > E > **–2.2359**	–0.0763 > E > **–2.6390**

TABLE II: d_s/δ for ERROR > ±5 % (HIGHER = BETTER)

Simulation	Wojda	Ferreira	Modified Wojda	Proposed Ferreira
Layers	**1.7310** < $\frac{d_s}{\delta}$ < 1.9463	**0.4414** < $\frac{d_s}{\delta}$ < 2.8633	**2.7335** < $\frac{d_s}{\delta}$ < 4.0525	**2.7474** < $\frac{d_s}{\delta}$ < 6.9914
Strands	**1.7632** < $\frac{d_s}{\delta}$ < 1.7867	**2.7845** < $\frac{d_s}{\delta}$ < 2.9007	**2.1348** < $\frac{d_s}{\delta}$ < 5.2156	**3.4153** < $\frac{d_s}{\delta}$ < 3.6081
Packing	**1.5930** < $\frac{d_s}{\delta}$ < 2.1726	**2.7050** < $\frac{d_s}{\delta}$ < 3.0388	**2.2590** < $\frac{d_s}{\delta}$ < 4.2533	**3.3260** < $\frac{d_s}{\delta}$ < 4.0733
Porosity	**1.7710**< $\frac{d_s}{\delta}$ < 4.7247	**1.4321** < $\frac{d_s}{\delta}$ < 2.8633	**2.1324** < $\frac{d_s}{\delta}$ < 5.1992	**2.8918** < $\frac{d_s}{\delta}$ < 3.6081

V. Experimental Validation

To validate the proposed methodology, experimental setups of short-circuited secondary transformers are prepared and measured using a Bode100. The winding structures and Litz wires used for the four prototypes are shown in Fig. 11. Four prototypes are prepared for the experimental validation according to the parameters from TABLE III.

TABLE III: PARAMETERS OF THE TESTED PROTOTYPES

Strand diameter d_s	Number of strands n_s	Turns per layer N_b	Layers per winding m
0.2 mm	400	6	1
0.1 mm	1000	7	1
0.1 mm	400	13	2
0.1 mm	150	20	2

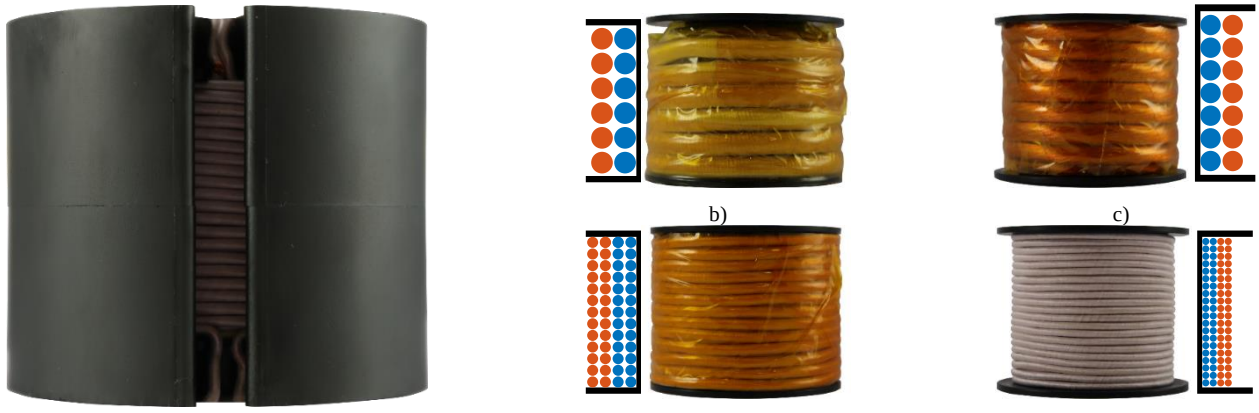

Fig. 11: Photograph of a) the assembled transformer for the resistance measurements and used winding structures. Litz wires used are b) 400x0.2, c) 1000x0.1, d) 400x0.1 and e) 150x0.1.

To avoid problems of having inner and outer core windings sections, common to EE core structures, a pot core that envelopes most of the windings is used. This ensures that the prototypes simulate the hypothesis used to generate the different 1-D models. The windings are excited using a 0.5 V sinusoidal signal. The impedance of the windings structure (primary connected to the Bode100, secondary short-circuited) is measured using the configuration from Fig. 12, and the real component used to extract the high frequency resistance.

The obtained results are shown in Fig. 13. Since all the models analyzed achieve similar accuracies at low and mid frequencies, the experimental results focus in the high and very high frequency behaviors ($1 < \Delta < 10$). Prototypes 3 (400x0.1) and 4 (150x0.1) (Fig. 11d and Fig. 11e) are made from multiple layers, resulting in high parasitic capacitances, thus the very high frequency range cannot be evaluated correctly. For prototypes 1 (400x0.2) and 2 (1000x0.1) (Fig. 11b and Fig. 11c), the parasitic capacitances are much lower, thus the transition from high to very high frequency can be seen clearly.

VI. Conclusions

In this work an in-depth analysis of 1-D models for the evaluation of high frequency winding losses is presented. First, a review of the existing approaches for individual wires is presented, demonstrating the basis of the different approaches based on Dowell's and Ferreira's approach. This is followed by an analysis of existing approaches to model Litz wire losses with 1-D models.

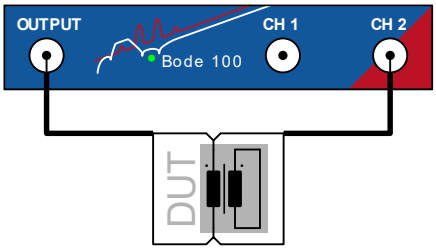

Fig. 12: Experimental setup using the Bode100 shunt-thru measurement.

Next, an alternative approach based on Ferreira's approach to model solid wire losses is presented. Instead of evaluating an average external field to all strands, the evolution of the external field per strand is modeled and used in the calculation of the losses, taking the form (32). This expression allows to evaluate the losses of each strand separately but can also be rewritten to (33) to obtain an average solution analogue to the rest of the models presented. Alternatively, a small modification of Wojda's model is also presented, slightly improving the obtained results.

Comparing against FEM simulation, the different models are evaluated and their sensibility to changes in the number of layers, number of strands, Litz packing factor and fill factor is presented. The results show how all models accurately estimate the low frequency losses, except Ferreira's approach in certain conditions. For high to very high frequency ranges, only Ferreira's approach and the proposed model based on Ferreira's approach are capable to attain accurate results. The proposed modification of Wojda's model is on average more accurate than the original model but is more sensible to all the analyzed variables and can overestimate the losses in extreme cases. Overall, the proposed model based on Ferreira's approach is

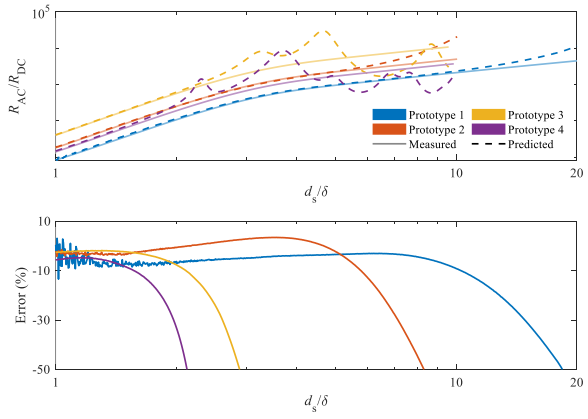

Fig. 13: Experimental measurements and predictions using the proposed Ferreira's approach-based model (32).

capable to accurately model the loses in a much wider frequency range. The capabilities of the proposed model to estimate the high frequency losses per strand are also demonstrated.

Lastly, four experimental prototypes using pot cores to achieve magnetic fields of only inner core sections are analyzed. It is shown that the proposed model correctly models the medium to high frequency ranges, but the parasitic capacitances of the multilayer prototypes does not allow to evaluate the losses at very high frequencies. On the contrary, the high to very high frequency transition is visible in the single layer prototypes, demonstrating the capabilities of the proposed model.

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