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Measurement of Optical Properties of Aqueous Lithium Bromide Solutions with and without Surfactant for Potential Use in Interferometric Analysis of Absorption Systems

Irene Pérez de Luco, Aliaksandr Mialdun,* Peru Arroiabe,* Ane Errarte,* Valentina Shevtsova,* and M. Mounir Bou-Ali*




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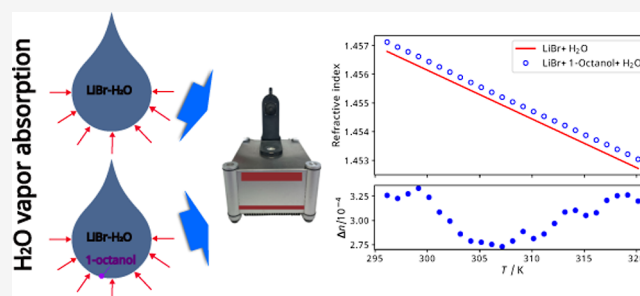
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ABSTRACT: The refractive indices of aqueous lithium bromide (LiBr–H₂O) solutions were measured over a wide range of mass fractions, from 0 to 0.62 kg/kg (62%), at seven wavelengths between 436.1 and 655 nm. Measurements were conducted over a temperature range of 296.15 to 320.15 K, with 1 K increments, using a commercial refractometer. In total, 3485 refractive index measurements were recorded, covering 35 compositions at 25 temperatures for each of the seven wavelengths. A parametrization model was proposed to unify the data set. The wavelength dependence of the refractive indices was further analyzed by using the Cauchy dispersion relation, from which three main coefficients were determined to predict $n(\lambda)$ for arbitrary wavelengths. Additionally, the optical contrast factors $\partial n/\partial T$ and $\partial n/\partial w$ were determined from the parametrized data. The effect of the surfactant (1-octanol) on the refractive index and contrast factors was carefully analyzed.



INTRODUCTION

Understanding the heat and mass transfer between absorbent solutions and vapor is critical for absorption refrigeration,¹ waste treatment,² air dehumidification, and pollutant absorption systems used in air and liquid purification, where absorbents trap pollutants.^{3,4}

Aqueous lithium bromide solutions (LiBr–H₂O) are among the best choices as working fluids in absorption chillers, heat transformers, and refrigerating systems.^{1,5} To achieve maximum efficiency, lithium bromide absorption chillers need to operate at a certain temperature and composition. It is therefore essential to determine the concentration of the solution at various locations as a function of temperature and time. A reliable and convenient method for measuring the concentration is through the refractive index, which can be obtained by either withdrawing a sample or, more invasively, in situ measurements within the system.

Advances in lasers and image processing techniques over the last few decades have enabled digital interferometry to drastically increase the accuracy of measurements. For monitoring concentration and temperature changes in transparent media, the interferometric approach is increasingly in demand^{6,7} as it effectively correlates with their temporal changes. In absorption research, this is particularly useful for tracking system dynamics in the presence of a surfactant. The application of interferometric methods provides valuable insight into the performance and efficiency of a system but

requires knowledge of the thermal and solutal contrast factors, which represent the derivatives of the refractive index with respect to temperature and concentration. Accurately measured contrast factors are essential for developing experimental techniques to study convection and heat and mass transfer in absorption and diffusive systems.

Although in situ composition tracking can be challenging, complementary numerical simulations^{8–10} are highly useful in understanding the flow dynamics. For a meaningful comparison with the experimental data, it is often necessary to recalculate the numerical results into the optical phase, which also requires knowledge of the optical contrast factors.

The available data on the refractive index of LiBr are limited in both temperature and composition ranges.^{11,12} In the range of compositions of practical interest, there appears to be a discrepancy in the dynamics of $n(w)$ behavior in these two sources. Here w is the mass fraction of LiBr in the binary mixture LiBr–H₂O. Furthermore, the existing data were obtained at a single wavelength, corresponding to the sodium

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Table 1. Chemical Specifications as per the Suppliers

IUPAC name	abbreviation	molecular weights	CAS Reg. No	suppliers	purity	purification
Lithium bromide	LiBr	86.85	7550–35–8	Sigma-Aldrich Company	≥0.99 kg/kg	None
1-Octanol	CH ₃ (CH ₂) ₇ OH	130.23	111–87–5	Sigma-Aldrich Company	≥0.99 kg/kg	None
Deionized Water	H ₂ O	18.02	7732–18–5	Autwomatic Plus system of Wasserlab	conductivity <1 μS/cm	none

66 D line, $m_D = 589.3$ nm, which may not provide sufficient
67 accuracy for comprehensive analysis across varying conditions.

68 This paper reports the refractive indices of LiBr solutions as
69 a function of mass fraction up to 0.62 kg/kg (62%) over a
70 temperature range of 296.15 to 320.15 K. Furthermore,
71 refractive indices were measured at seven wavelengths. Optical
72 contrast factors, $\partial n/\partial T$ and $\partial n/\partial w$, were calculated from the
73 parametrized data. We also examined the Cauchy dispersion
74 relation. In addition, the refractive indices of LiBr solutions in
75 the presence of the surfactant octanol were measured. While
76 small quantities of surfactants generally have a minimal impact
77 on many physical properties, they can affect optical properties.
78 The specific objective of this study is to provide data that will
79 improve our understanding of the behavior of absorbent
80 materials. These results significantly extend the database of this
81 important absorbent.

82 ■ EXPERIMENTAL METHODS

83 **Solutions Preparation.** The chemicals used in this study
84 together with their basic information are summarized in Table
85 1. Given the extreme hygroscopic nature of lithium bromide, it
86 is essential to follow a strict procedure to prevent unintentional
87 absorption of moisture from the atmosphere during sample
88 preparation. Lithium bromide (LiBr), with a mass fraction
89 purity of ≥99% (CAS 7550–35–8), was obtained from Sigma-
90 Aldrich Company. Before the samples were prepared, the LiBr
91 was heated in an oven at 120 °C for a minimum of 16 h.
92 Immediately after removal from the oven, LiBr was weighed
93 directly into the sealed jar used for sample preparation to
94 minimize exposure to atmospheric moisture.

95 A digital mass balance Gram VXI-310 was employed to
96 measure both the LiBr and water masses with a combined
97 standard uncertainty of ±0.0001 g. All solutions were prepared
98 as mass fractions w [kg/kg] of anhydrous LiBr.

99 We prepared 35 solutions of different compositions covering
100 the composition range from 0 up to $w = 0.62$ kg/kg.

101 **Refractive Index Measurements.** The data on refractive
102 index, as well as the derived solutal and thermal contrast
103 factors, were obtained using a multiwavelength Anton Paar
104 Abbemat WR-MW refractometer. The apparatus was cali-
105 brated using distilled deionized water. Calibration was
106 performed daily prior to measurements at the reference
107 temperature of 298 K. The instrument's internal reference
108 data, which are based on measurements relative to air, were
109 used for this purpose. We verified that these internal values are
110 consistent with the IAPWS formulation,¹³ appropriately
111 adjusted to the air reference frame. For each reported value,
112 three individual measurements were collected and averaged.
113 Before each measurement with a fresh solution, the apparatus
114 was thoroughly cleaned with water and acetone. The mean
115 values and standard deviations were calculated. The temper-
116 ature of the sample and the internal temperature of the optical
117 system were maintained within ±0.03 K due to the two built-in
118 Peltier devices. The instrument operates at seven different

wavelengths of visible light, covering the range of $\lambda = 436$ –655
nm, ensuring a precision of $\pm 4 \times 10^{-5}$ nm.

121 **Comparison with Available Results.** Figure 1 presents a
122 comparison between our experimental data at $\lambda = 589.3$ nm

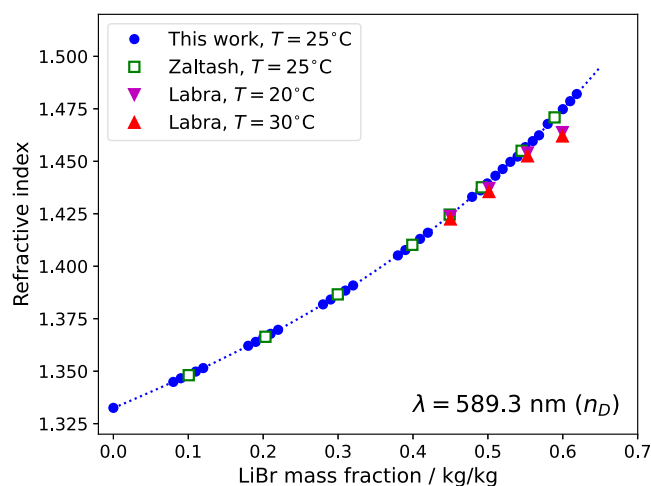


Figure 1. Comparison of measured data with literature values. This work - blue circles, values reported by Zaltash¹¹ - open green squares; values reported by Labra et al.¹² are given by magenta and red triangles. The dotted interpolation curve is for guidance only.

and the literature data reported by Zaltash¹¹ and Labra et al.¹²
Our and Zaltash's data are taken at $T = 298.18$ K; Labra's data
are not available at this temperature; instead, values for $T =$
293.15 K and $T = 303.15$ K are given. The comparison is given
in terms of the dependence $n(w)$ and is in excellent agreement
with the Zaltash data and consistent with the data of Labra et
al..

■ RESULTS AND DISCUSSION

Refractive Index of LiBr–H₂O Solutions. In this work,
the experimental data of the refractive index (n) for the LiBr–
water binary mixture were measured over the whole range of
composition up to the saturation point, at temperatures from $T =$
296.15 K up to $T = 320.15$ K with a 1 K increment at 35
different compositions. The seven tables with raw data are
given in Supporting Information.

Figure 2a presents the variation of the refractive index with
temperature measured at different wavelengths of the
composition of practical interest, $w = 0.55$ kg/kg. The general
trend is a linear dependence when n decreases with an increase
in temperature, which can be correlated using the following
equation

$$n_{\lambda}(T) = a_{\lambda}^{00} + a_{\lambda}^{10} \times (T - T_0) \quad (1)$$

where a_{λ}^{00} and a_{λ}^{10} are adjustable parameters depending on
wavelength and $T_0 = 298.15$ K. The slopes a_{λ}^{10} , which are
associated with optical contrast factor $\partial n/\partial T$, weakly decrease
from $a_{\lambda}^{10} = -1.73 \cdot 10^{-4}$ to $a_{\lambda}^{10} = -1.69 \cdot 10^{-4}$ as the wavelength

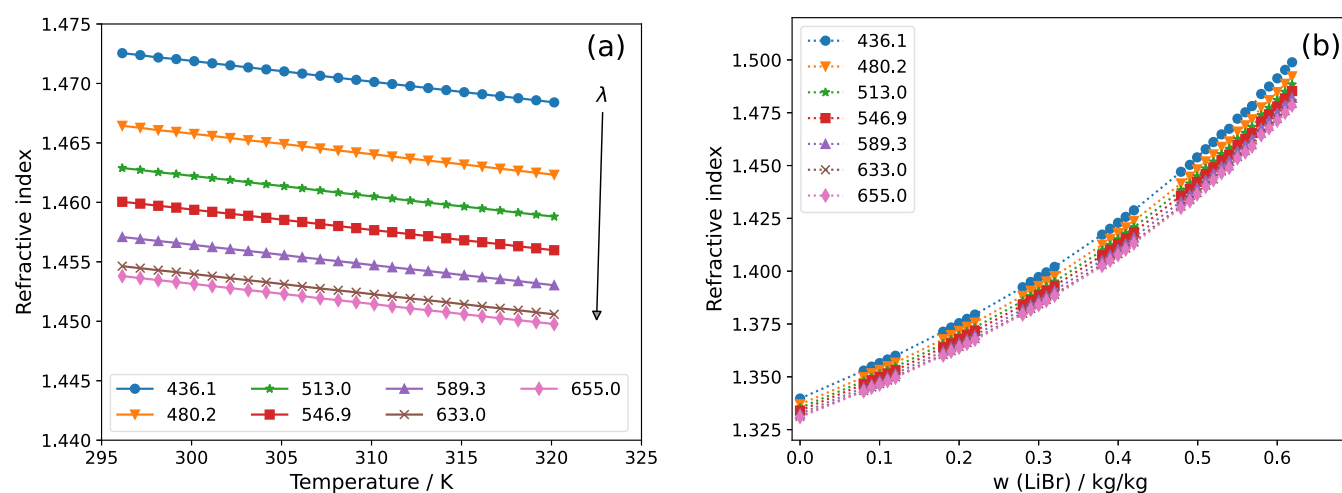


Figure 2. Experimental refractive indices. (a) Variation of refractive index with temperature measured at different wavelengths at $w = 0.55$ [kg/kg]; (b) variation of refractive index with mass fraction of LiBr measured at different wavelengths at $T = 298.15$ K. The curves are for guidance only.

149 varies from $\lambda = 436$ nm to $\lambda = 655$ nm. Note that these values
150 are consistent only for a mass fraction $w = 0.55$ kg/kg.

151 Panel (b) shows the variation of the refractive index with the
152 mass fraction of LiBr measured at different wavelengths at $T =$
153 298.15 K. It can be seen that the 35 compositions studied are
154 divided into seven groups for convenience. The last three
155 groups include continuous measurements, as this area is of
156 particular interest for chillers. The raw data for one
157 composition in each group, i.e., $w = 0.1, 0.2, 0.3, 0.4, 0.5,$
158 0.55, and 0.6 kg/kg are given as a single table (Tables S1–S7
159 in the Supporting Information) including the n values at 25
160 different temperatures and at seven wavelengths.

161 Unlike temperature, the composition dependence on the
162 refractive index exhibits a third order polynomial behavior,
163 which requires the determination of four adjustable coefficients
164 for each temperature; see details in the next section. As can be
165 seen from Figure 3, the refractive index decreases with
166 wavelength. However, the effect is quite weak.

167 To extend the measured data to arbitrary wavelengths, we
168 examine the dispersion relations between refractive index n and
169 wavelength λ . Aqueous solutions are polar and inherit the

strong absorption features of water in the infrared.^{14,15} Our
170 measurements include deep UV or IR and are made at a
171 wavelength range of $455 \text{ nm} \leq \lambda \leq 655 \text{ nm}$. Among the
172 various dispersion equations, the Sellmeier and the Cauchy
173 formula are the most common.¹⁵ Careful analysis of molar
174 refractivity of water and/or dilute water vapor in literature
175 provides the basis that the use of Cauchy equation in the form
176

$$n(\lambda) = A + \frac{B}{\lambda^2} + \frac{C}{\lambda^4} + \dots \quad (2)$$

is justified in the considered range of wavelengths. Here $A,$
178 $B,$ and C are coefficients that can be determined by fitting them
179 to the measured refractive indices.

We have calculated all three coefficients, and they can be
181 found in the Supporting Information in the last columns of all
182 seven tables.

Our previous studies of nonpolar mixtures^{17,18} showed that
184 adding the $1/\lambda^4$ to the Cauchy equation improves the accuracy
185 of the interpolation by an order of magnitude with resulting
186 MAPD (mean absolute percentage deviation) on the level of
187 0.002%.

$$\text{MAPD} = \frac{1}{N} \sum_{i=0}^N \left| \frac{n_i^{\text{exp}} - n_i^{\text{fit}}}{n_i^{\text{fit}}} \right| \times 100\% \quad (3)$$

Parameterization of the Measured $n_\lambda(w, T)$ Values.

The temperature and concentration dependencies of refractive
191 index $n_\lambda(T, w)$ were parametrized independently for each
192 wavelength. The parametrization is done with the use of
193 statsmodels library for Python programming language.¹⁹
194 The raw data at all measured temperatures and concentrations
195 were fitted to a polynomial equation of the following form
196

$$n_\lambda(T', w) = \sum_{i=0}^N \sum_{j=0}^M a_\lambda^{(ij)} T'^i w^j \quad (4)$$

where $T' = T - T_0$ is the temperature relative to the reference
198 $T_0 = 298.15$ K. The highest polynomial orders N and M with
199 respect to temperature and concentration, respectively, were
200 found by analyzing the corresponding dependencies individ-
201 ually, using subsets with a constant second parameter. An
202 example of the analysis is shown in Figure 4, where the plotted
203 residues are calculated as

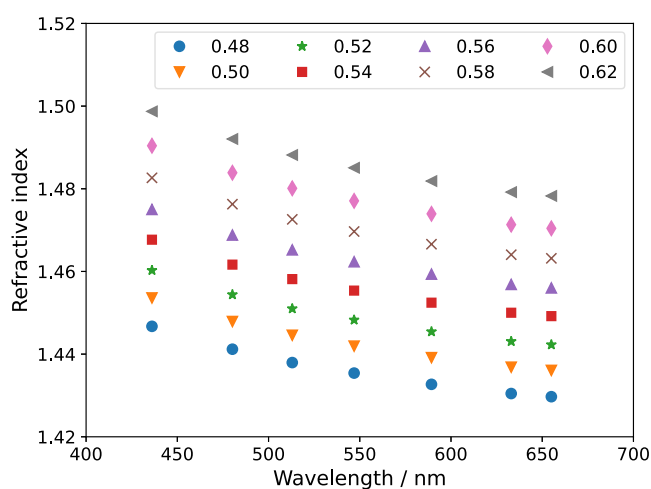


Figure 3. Variation of the refractive index with wavelength at $T = 303.15$ K for different mixture compositions [kg/kg], as indicated in the plot.

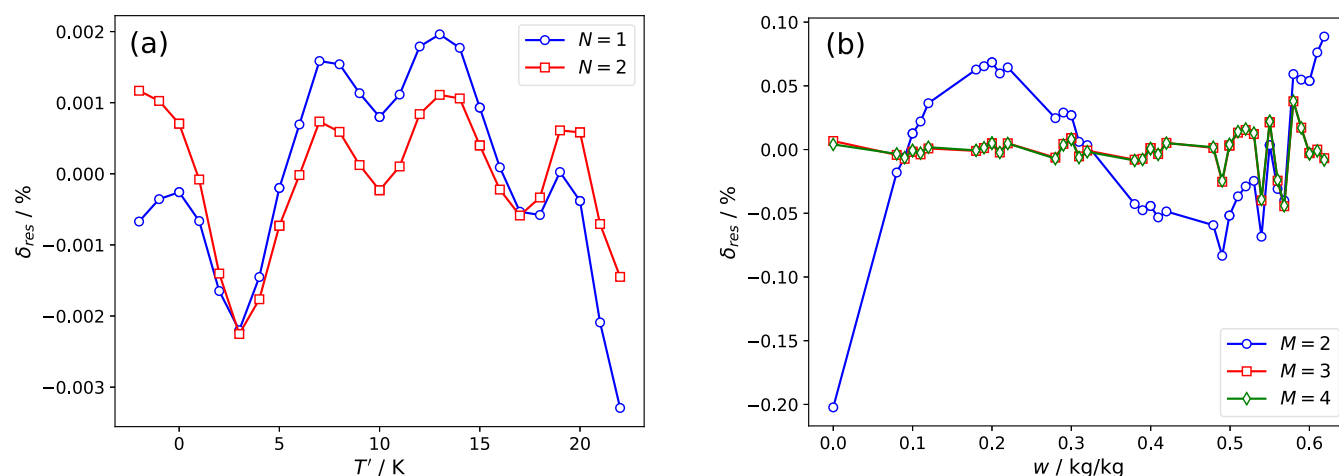


Figure 4. Selecting the highest orders of the polynomials N and M from a comparison of the residuals between the measured values and the fit to the data when either w or T are constant. (a) Temperature dependence of refractive index at 589.3 nm and $w = 0.30$ kg/kg. (b) Concentration dependence of refractive index at 589.3 nm and $T = 308.15$ K.

$$\delta_{\text{res}} = \frac{n_{\text{meas}} - n_{\text{fit}}}{n_{\text{fit}}} \times 100\% \quad (5)$$

205

As seen in Figure 2a, the temperature dependence of n_{λ} is close to linear. We tested the adequacy of the first-order polynomial by comparing it with the polynomial of the second-order (see Figure 4a for the results). This comparison demonstrates that the second order polynomial fitting does not provide an essential improvement in the accuracy; thus, the optimal polynomial order for the temperature dependence may be limited to $N = 1$.

It follows from Figure 2b that the dependence of n_{λ} on the concentration is nonlinear. To determine the most suitable polynomial fit, we tested polynomials of second and higher orders (see the comparison of the fits in Figure 4b). The polynomial of the second order is insufficient, while the polynomial of the fourth order does not improve the fit accuracy compared to the third order. Consequently, the optimal polynomial order for the concentration dependence was chosen as $M = 3$.

As a result, the final form of the data parametrization becomes

$$n_{\lambda}(T', w) = a_{\lambda}^{(00)} + a_{\lambda}^{(01)}w + a_{\lambda}^{(02)}w^2 + a_{\lambda}^{(03)}w^3 + a_{\lambda}^{(10)}T' + a_{\lambda}^{(11)}T'w + a_{\lambda}^{(12)}T'w^2 + a_{\lambda}^{(13)}T'w^3 \quad (6)$$

225

Alternatively, the eq 6 can be presented in the matrix form as

$$n_{\lambda}(T', w) = (1T') \times \begin{bmatrix} a_{\lambda}^{(00)} & a_{\lambda}^{(01)} & a_{\lambda}^{(02)} & a_{\lambda}^{(03)} \\ a_{\lambda}^{(10)} & a_{\lambda}^{(11)} & a_{\lambda}^{(12)} & a_{\lambda}^{(13)} \end{bmatrix} \begin{pmatrix} 1 \\ w \\ w^2 \\ w^3 \end{pmatrix} \quad (7)$$

227

Next, we examined a need in the higher order terms of eq 6, namely, the coefficients $a_{\lambda}^{(11)}$, $a_{\lambda}^{(12)}$, and $a_{\lambda}^{(13)}$. This is an essential step, as the presence of excessive terms may increase the uncertainty of the high-order coefficients. This analysis was done by comparing the fit results with all eight coefficients $a_{\lambda}^{(ij)}$ (hereafter 8-parameter fit), with removed $a_{\lambda}^{(13)}$ coefficient (7-parameter fit), with removed $a_{\lambda}^{(12)}$ and $a_{\lambda}^{(13)}$ coefficients (6-

parameter fit), and with removed all three high-order coefficients $a_{\lambda}^{(11)}$, $a_{\lambda}^{(12)}$, $a_{\lambda}^{(13)}$ (5-parameter fit). The visualization and comparison of the fit results are convenient via the concentration dependence of the temperature contrast of the refractive index, $\partial n/\partial T$, presented in Figure 5.

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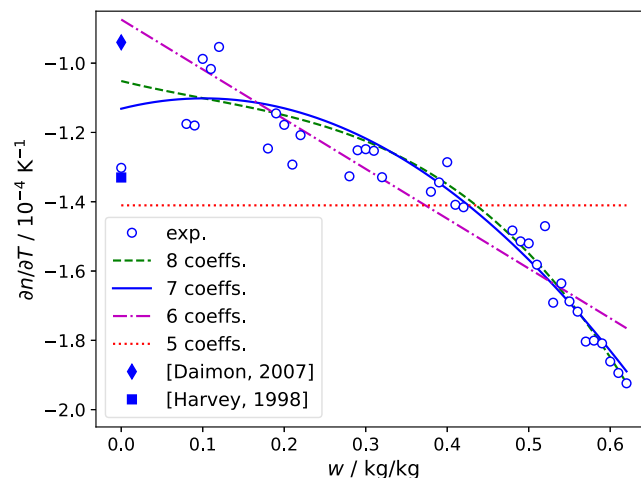


Figure 5. Comparison of the fit results with different numbers of high-order terms in the parametrization by eq 6 for the refractive index at 633 nm and $T = 308.15$ K. The literature data on the contrast factor of pure water are taken from ref 13 at 633 nm and 308.15 K and from ref 14 at 644 nm and 294.65 K.

The results of the different fits are presented in the plot by curves, while the experimental points are obtained by independent linear interpolation of the measurements at all assessed w_i .

243

Obviously, limiting the fit to 5 coefficients is not a good option, as it leads to $\partial n/\partial T = \text{constant}$ in the full concentration range. The fit with 6 coefficients provides a linear dependence $\partial n/\partial T = f(w)$, which is better but still does not outline the contrast factor behavior in diluted solutions. The 7-parameter fit bends the contrast factor dependence toward the literature data for pure water from refs. 13,14, which is also consistent with our own data. The dependence $n(T)$ for pure water is nonlinear in a wide temperature range; thus, for comparison with our results we selected the values specific to $T = 308.15$

253

Table 2. Coefficients of the Polynomial Parameterization of the Refractive Index of LiBr–Water Solutions in the Concentration Range of 0.0–0.60 Mass Fraction and the Temperature Range of 296.15 < T < 320.15 K at Seven Measured Wavelengths

Coeff	wavelength λ , nm						
	436.1	480.2	513.0	546.9	589.3	633.0	655.0
$a_{\lambda}^{(00)}$	1.33971 (0.00009)	1.33692 (0.00009)	1.33522 (0.00009)	1.33390 (0.00009)	1.33244 (0.00009)	1.33120 (0.00009)	1.33082 (0.00009)
$a_{\lambda}^{(01)}/10^{-1}$	1.6118 (0.0089)	1.5747 (0.0086)	1.5554 (0.0085)	1.5379 (0.0084)	1.5224 (0.0083)	1.5068 (0.0082)	1.5047 (0.0083)
$a_{\lambda}^{(02)}/10^{-1}$	0.499 (0.027)	0.460 (0.026)	0.433 (0.026)	0.419 (0.026)	0.394 (0.025)	0.386 (0.025)	0.369 (0.025)
$a_{\lambda}^{(03)}/10^{-1}$	1.709 (0.026)	1.704 (0.025)	1.706 (0.025)	1.700 (0.025)	1.707 (0.024)	1.700 (0.024)	1.710 (0.024)
$a_{\lambda}^{(10)}/10^{-4}$	-1.186 (0.062)	-1.170 (0.060)	-1.162 (0.060)	-1.141 (0.059)	-1.139 (0.058)	-1.132 (0.057)	-1.120 (0.058)
$a_{\lambda}^{(11)}/10^{-4}$	0.74 (0.41)	0.67 (0.40)	0.65 (0.40)	0.57 (0.39)	0.59 (0.39)	0.59 (0.38)	0.57 (0.39)
$a_{\lambda}^{(12)}/10^{-4}$	-3.15 (0.58)	-3.03 (0.57)	-3.00 (0.56)	-2.90 (0.55)	-2.92 (0.55)	-2.93 (0.54)	-2.93 (0.55)
$a_{\lambda}^{(13)}/10^{-4}$	0	0	0	0	0	0	0
MAPD, %	0.0095	0.0093	0.0091	0.0090	0.0089	0.0087	0.0087

254 K, which is a middle of the temperature interval explored in the
 255 present work. The full 8-parameter fit does not improve the
 256 agreement between the experiment and the fit but even
 257 worsens the agreement for the low salt concentration.
 258 Consequently, we selected a 7-parameter fit as a working
 259 approach, which in practice means dropping the last term from
 260 the RHS of eq 6 or setting $a_{\lambda}^{(13)} = 0$ in eq 7.

261 The coefficients for fitting the data at different wavelengths
 262 are summarized in Table 2.

263 **Refractive Index of LiBr–H₂O Solution in the**
 264 **Presence of Surfactant.** Surfactants are used to enhance
 265 the absorption process.²⁰ In the presence of even a small
 266 disturbance, they cause additional surface tension gradients in
 267 the LiBr–water solution. These gradients, in turn, promote
 268 liquid motion near the interface, a phenomenon referred to as
 269 Marangoni convection.^{9,10} Two common surfactants for LiBr
 270 aqueous solutions are 1-octanol and 2-ethyl-1-hexanol.

271 Hereafter, we denote the concentration of the surfactant as
 272 w_s . The optical properties of the LiBr–water system were
 273 studied with the addition of a small amount ($w_s \approx 140 - 170$
 274 ppm) of 1-octanol (CAS 111–87–5) to solutions of LiBr–
 275 H₂O in the range of compositions $w = 0.48 - 0.62$ kg/kg and at
 276 temperatures of 298.15, 303.15, 308.15, 313.15, and 318.15 K.
 277 This range of LiBr mass fractions is of most practical interest.
 278 The uncertainty of the mass fraction of the surfactant is 0.066.
 279 The experimental results are shown in Figure 6a in open
 280 circles. Measurements reveal that the addition of 1-octanol
 281 affects the refractive index of LiBr/H₂O solutions in a
 282 composition-dependent manner. The largest change in
 283 refractive index was observed at a LiBr mass fraction of $w =$
 284 0.50 kg/kg, where the mean deviation, $\Delta n \sim 4.2 \times 10^{-4}$, was
 285 consistent across all seven wavelengths, corresponding to a
 286 change of about 0.008%.

287 However, comparing measurements directly between
 288 systems with and without the surfactant can be misleading
 289 due to small differences in the concentration of the basic LiBr/
 290 H₂O solutions used in each set of measurements. To minimize
 291 the influence of these outlier points, where some of the
 292 variation in Δn may be due to the LiBr concentration deviation
 293 rather than the surfactant, the measurements with surfactant

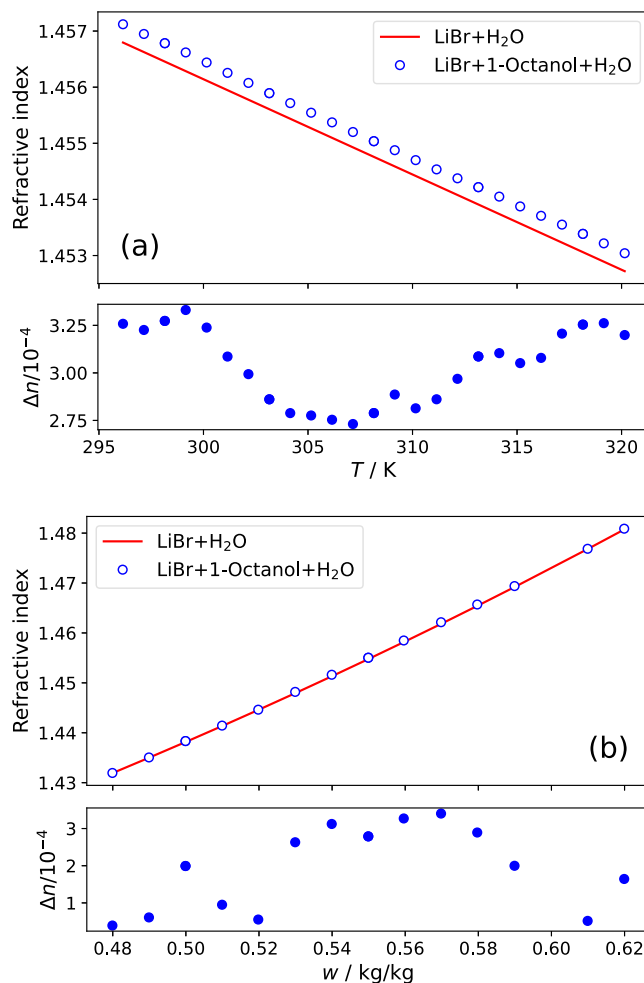


Figure 6. Refractive index in the LiBr–H₂O system with (1-octanol) surfactant and without (using eq 6) at $\lambda = 589.3$ nm and the difference Δn between them. (a) Temperature dependence at $w = 0.55$ kg/kg and (b) concentration dependence at $T = 308.15$ K.

294 are compared to the parametrized refractive index of the LiBr–
295 water solution without 1-octanol, as given by eq 6.

296 Figure 6 presents a comparison of the refractive index values
297 for LiBr–H₂O solutions with and without the surfactant, using
298 the parametrized n values, along with the corresponding Δn
299 caused by the surfactant at the bottom panel. This comparison
300 confirms our suspicion regarding the compositional variation,
301 as there is a small but noticeable difference between the
302 parametrized refractive index values for $w = 0.55$; see the upper
303 plot. This difference is not visible in Figure 6a. The
304 parametrized values are smoother, and they show a nearly
305 constant difference in refractive index between the solutions
306 with and without surfactant for all studied compositions.

307 The analysis of Figure 6 shows that the effect of adding 1-
308 octanol on the refractive index has a minor variation with
309 temperature (see panel a) and a slightly stronger effect with the
310 composition of the LiBr–water solutions (as seen in panel b).

311 Overall, excluding some outliers, across the covered range of
312 T and w the surfactant additive gives a consistent positive
313 deviation to the refractive index falling into the range $\Delta n =$
314 $0.5\text{--}3.3 \cdot 10^{-4}$.

315 **Optical Contrast Factors.** Contrast Factors of LiBr–H₂O
316 Solutions without Surfactant. The expressions for the
317 contrast factors can be obtained by differentiating eqs 6 or 7.
318 The term $a_\lambda^{(13)}w^3$ is omitted, as discussed above. The thermal
319 optical contrast factors are independent of temperature and
320 exhibit a quadratic dependence on composition

$$\left. \frac{\partial n_\lambda}{\partial T} \right|_w = a_\lambda^{(10)} + a_\lambda^{(11)}w + a_\lambda^{(12)}w^2 \quad (8)$$

322 The solutal optical contrast factors exhibit a linear dependence
323 on the temperature and a quadratic dependence on
324 composition

$$\left. \frac{\partial n_\lambda}{\partial w} \right|_T = a_\lambda^{(01)} + 2 \times a_\lambda^{(02)}w + 3 \times a_\lambda^{(03)}w^2 + a_\lambda^{(11)}T' + 2 \times a_\lambda^{(12)}T'w \quad (9)$$

326 The advantage of the matrix representation of the para-
327 metrization is that the same matrix of the coefficients

$$[a_\lambda^{(ij)}] = \begin{bmatrix} a_\lambda^{(00)} & a_\lambda^{(01)} & a_\lambda^{(02)} & a_\lambda^{(03)} \\ a_\lambda^{(10)} & a_\lambda^{(11)} & a_\lambda^{(12)} & 0 \end{bmatrix} \quad (10)$$

329 can be used both for refractive index and the contrast factors
330 calculation, with corresponding modification of the vectors of
331 powers for T' and w , as follows

$$\left. \frac{\partial n_\lambda}{\partial T} \right|_w = (01)[a_\lambda^{(ij)}] \begin{pmatrix} 1 \\ w \\ w^2 \\ w^3 \end{pmatrix} \quad (11)$$

333 and

$$\left. \frac{\partial n_\lambda}{\partial w} \right|_T = (1T')[a_\lambda^{(ij)}] \begin{pmatrix} 0 \\ 1 \\ 2w \\ 3w^2 \end{pmatrix} \quad (12)$$

The contrast factors of LiBr–H₂O solutions, required for
optical experiments, are shown in Figure 7 and are given in the

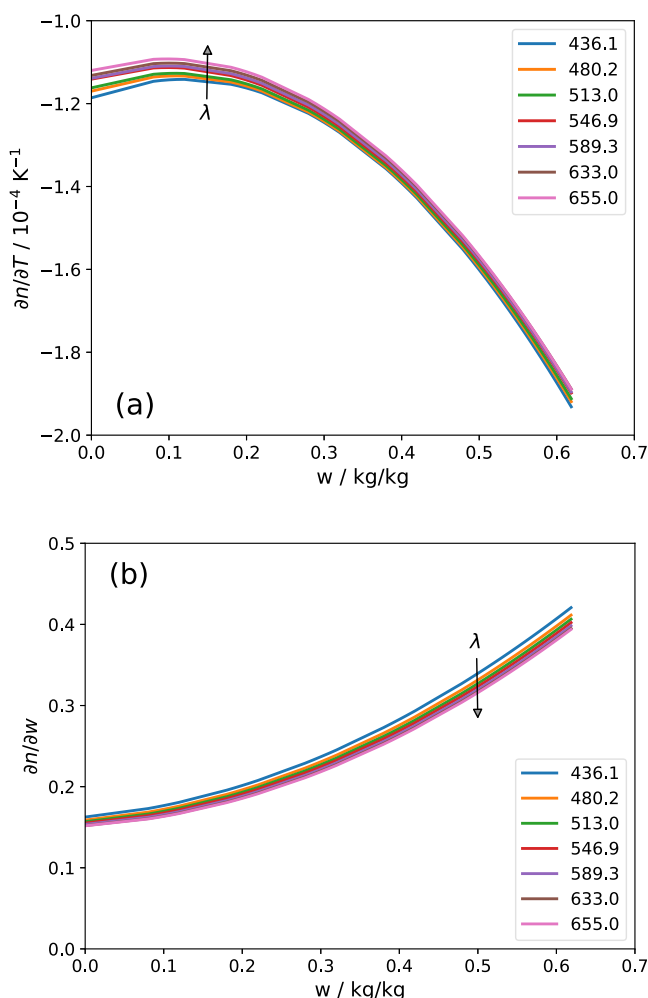


Figure 7. Variation of (a) thermal and (b) solutal contrast factors with LiBr mass fraction at seven wavelengths indicated in the plot at $T = 298.15$ K ($T' = 0$).

form of tables in the Supporting Information. They are
calculated using eqs 8 and 9 with the parametrization
coefficients listed in Table 1. Figure 7 illustrates that the
dependence of the contrast factor on the wavelength is
generally not significant. However, the thermal contrast factor
($\partial n/\partial T$) exhibits a larger dependence on λ for lower mass
fractions, while $\partial n/\partial T$ exhibits a larger dependence at higher
mass fractions. Another observation is that $\partial n/\partial T$ decreases
with w , whereas $\partial n/\partial w$ increases with w . It is important to note
that $\partial n/\partial T$ is of the order of 10^{-4} , while $\partial n/\partial w$ is of the order
of 1.

Contrast Factors of LiBr–H₂O Solutions with Surfactant
1-Octanol. Two approaches were applied to study how the
surfactant affects the optical properties of the mixture. First,
the influence of the surfactant on the thermal contrast factor
was examined. Already looking at Figure 7a, one may
anticipate a similar magnitude of $\partial n/\partial T$ for the mixtures
without and with surfactant. This is confirmed by Figure 8,
which summarizes the thermal contrast factor for both types of
mixtures as a function of salt concentration, w , supporting the
preliminary observation of their similarity.

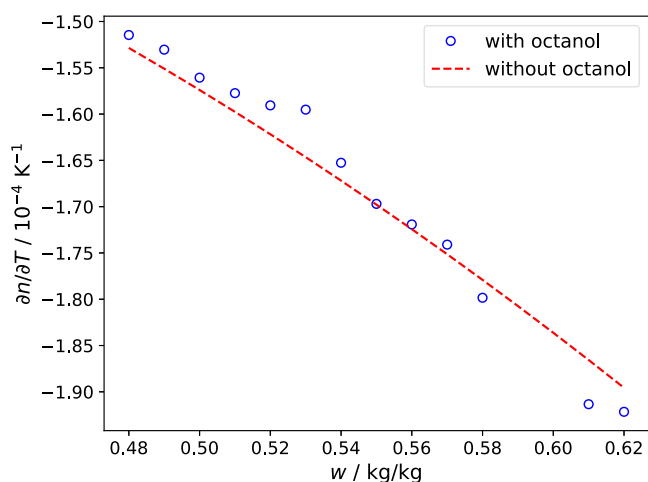


Figure 8. Variation of thermal contrast factor with LiBr mass fraction at 589.3 nm wavelength for the mixtures without and with surfactant.

358 Next, we sought to estimate the contrast factor of 1-octanol
 359 at infinite dilution as well as its potential dependence on
 360 temperature and concentration. The estimation of $\partial n/\partial w_s$ was
 361 conducted as follows: the refractive index excess due to the
 362 surfactant was calculated as the difference between the
 363 measured value for the mixture with surfactant and the
 364 parametrized value for the mixture without octanol,
 365 $\Delta n = n_{w+w_s}^{\text{exp}} - n_w^{\text{param}}$. The contrast factor was then determined
 366 by dividing the refractive index excess by the accurately
 367 measured concentration of 1-octanol.

$$\left. \frac{\partial n}{\partial w_s} \right|_{T,w} = \frac{\Delta n}{w_s} \quad (13)$$

369 While we did not parametrize the contrast factor $\partial n/\partial w_s$ over
 370 all assessed range of parameters, we estimated it for selected
 371 cases to gain a general understanding of its magnitude. The
 372 results are presented in Figure 9.

373 From the plot in Figure 9a, it can be observed that the
 374 temperature dependence of $\partial n/\partial w_s$ is very weak if present at
 375 all. In contrast, the concentration of basic salt w has a
 376 significant, nonlinear effect on the contrast factor, as shown in
 377 Figure 9b.

378 For the most relevant concentrations, the contrast factor at
 379 infinite dilution of w_s falls within the range of $\partial n/\partial w_s = 1 - 2$,
 380 which is 2 to 4 times higher than the maximum solutal contrast
 381 factor for the LiBr–H₂O solutions. This pronounced optical
 382 effect of 1-octanol is expected due to its large and heavy
 383 molecular structure. However, the measured contrast factor for
 384 octanol is still several times lower compared to that of polymer
 385 solutions.

386 ■ CONCLUSIONS

387 The study provides a comprehensive data set of 6125 refractive
 388 index measurements for aqueous LiBr solutions across a wide
 389 range of compositions, temperatures, and wavelengths. In
 390 addition to the measured values of n at seven wavelengths, the
 391 coefficients for the Cauchy dispersion relation were carefully
 392 determined, offering a reliable method to predict the refractive
 393 index for arbitrary wavelengths. The measured refractive index
 394 data were parametrized, improving the utility of the data and
 395 minimizing the impact of outliers.

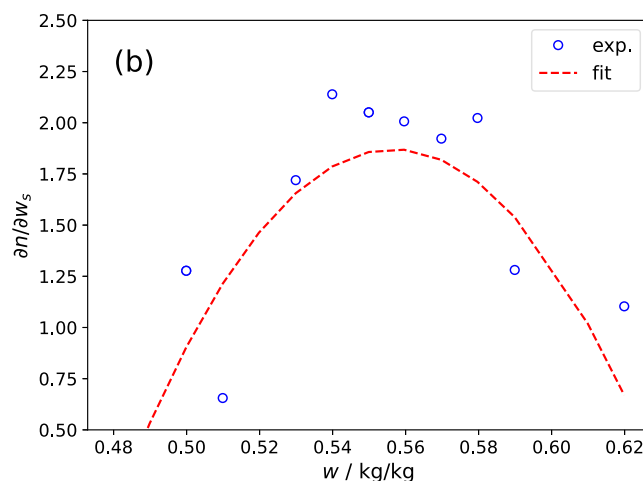
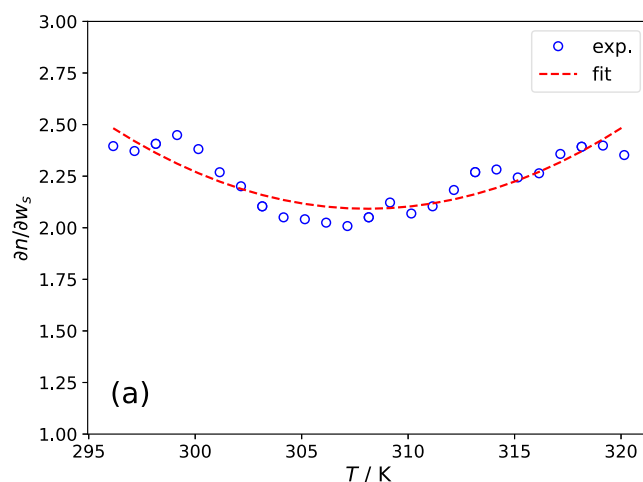


Figure 9. Solutal contrast factor of 1-octanol at infinite dilution in the LiBr–H₂O mixture at $\lambda = 589.3$ nm. (a) Temperature dependence at $w = 0.55$ kg/kg and (b) concentration dependence at $T = 308.15$ K.

The optical contrast coefficients, $\partial n/\partial T$ and $\partial n/\partial w$,
 396 calculated from the parametrized data provide important
 397 information about the temperature and concentration depend-
 398 encies of the refractive index. This may be of importance for
 399 optical and thermophysical applications, such as heat transfer
 400 and concentration gradient processes involving LiBr solutions.

The effect of the surfactant (1-octanol) on the refractive
 402 index and contrast factors was also carefully analyzed.
 403 Additionally, we determined the contrast factors of 1-octanol
 404 at an infinite dilution.

406 ■ ASSOCIATED CONTENT

407 Supporting Information

The Supporting Information is available free of charge at
 408 <https://pubs.acs.org/doi/10.1021/acs.jced.5c00066>.

Thermal ($\partial n/\partial T$) and solutal ($\partial n/\partial w$) contrast factors
 410 of LiBr–H₂O solutions at $T = 298$ K; refractive indices
 411 of LiBr–H₂O solutions and Cauchy coefficients for
 412 mass fractions $w = 0.1, 0.2, 0.3, 0.4, 0.5, 0.55,$ and 0.60
 413 kg/kg at temperature range $295 \text{ K} \leq T \leq 320 \text{ K}$; and
 414 refractive indices of LiBr–H₂O with the addition of
 415 surfactant, 150 ppm of 1-octanol, at temperature range
 416 $295 \text{ K} \leq T \leq 320 \text{ K}$ (PDF)
 417 All measured data (3485 points) (XLSX) 418

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448 Notes

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